KdV soliton solutions to a model of hepatitis C virus evolution

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Russell and the Wave of Translation

"<...> the boat suddenly stopped - not so the mass of water in the channel which it. had put in motion; it accumulated round the prow of the vessel in a state of violent agitation, then suddenly leaving it behind, rolled forward with great velocity, assuming the form of a large solitary elevation <...> which continued its course along the channel apparently without change of form or diminution of speed. <...> was my first chance interview with that singular and beautiful phenomenon which I have called the Wave of Translation. (1834)"

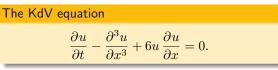


Figure: John Scott Russell (1808 – 1882)

First Mathematical Model: KdV equation



Figure: D. Korteweg (top), G. de Vries The first actual mathematical model of solitary waves (solitons) was discovered by **Boussinesq** (1877) and later rediscovered and studied in detail by **Diederik Korteweg** and **Gustav de Vries**.



It was shown that there exist solutions of the form:

KdV soliton

$$u(t,x) = -\frac{c}{2}\operatorname{sech}^{2}\left(\frac{\sqrt{c}}{2}\left(x - ct - a\right)\right).$$

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Zabusky and Kruskal: Rediscovery

The works of Russell, Boussinesq, Korteweg and de Vries fell into obscurity until 1965, when **Norman J. Zabusky** and **Martin Kruskal** made connections between the KdV equation and the Fermi-Pasta-Ulam experiment.

The word "soliton" was coined and extensive studies into the nature of soliton (solitary) processes were launched that are still continuing today. It has had a broad and far-reaching impact in myriad fields ranging from the purest mathematics to experimental science.



Figure: N. Zabusky (top), M. Kruskal

Rise to fame: the Schrödinger equation

A particular surge in interest of the analysis of solitary processes in physics came when **Vladimir Zakharov** and **Aleksei Shabat** demonstrated in 1972 that the nonlinear Schrödinger (NLS) equation has soliton solutions.

NLS equation $i \frac{\partial u}{\partial x} + \frac{\partial^2 u}{\partial t^2} \pm 2 |u|^2 u = 0.$

This discovery had enormous repercussions in physics, especially in nonlinear optics, Bose-Einstein condensates, where the NLS equation plays a very important role.



Figure: V. Zakharov (top), A. Shabat

Applications of Soliton Theory

Soliton theory has had a great impact on myriad fields of science, including:

- Nonlinear optics;
- Bose-Einstein condensates;
- Hydrodynamics;
- Biophysics;
- MEMs and NEMs;
- Plasmas;
- Population dynamics.

Physical Properties of Solitary Solutions

A nonlinear wave is called a soliton if:

- It maintains its shape as it propagates at a constant speed;
- If it collides with another soliton, it emerges from the collision unaltered, except for a phase shift.

The definition is not universally accepted – there are a few ways to define solitons in the physical sense (for example allowing a small loss of energy after collision ("light bullets")), however, from a mathematical perspective the definition given above is almost ubiquitous.

Solitary solutions - analytic expression

Solitary solution

The m-th order solitary solution reads:

$$x(t) = \frac{\sum_{k=0}^{m} \alpha_k \exp\left(\eta k(t-c)\right)}{\prod_{k=1}^{m} \left(\exp\left(\eta(t-c)\right) - t_k\right)},$$

where $\eta, \alpha_k, t_k, c \in \mathbb{C}$ are fixed parameters.

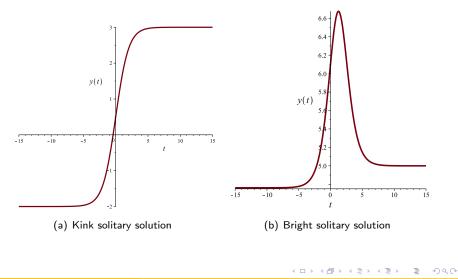
The shape is based on the **KdV soliton**, which is a special case of the above solution (hyperbolic secant):

$$\varphi\left(\xi\right) = -\frac{c}{2}\operatorname{sech}^{2}\left(\frac{\sqrt{c}}{2}\left(\xi - \xi_{0}\right)\right).$$

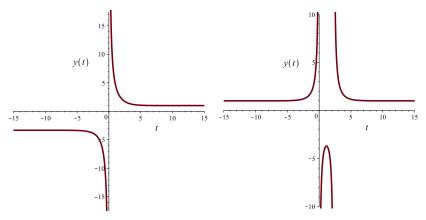
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Examples

Solitary solutions



Solitary solutions



(c) Solitary solution with one singularity

(d) Solitary solution with two singularities

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Coupled Riccati equations (1)

Regula et al (2009) introduced the following system for modeling Hepatitis C virus (HCV) evolution:

$$\begin{split} & x'_t = x \left(1 - x - y \right) - (1 - \theta) \, bxy + qy + s; \\ & y'_t = ry \left(1 - x - y \right) + (1 + \theta) \, bxy - (d + q)y, \\ & \theta, b, q, s, r, d \in \mathbb{R}. \end{split}$$

Simple generalization leads to:

$$\begin{aligned} x'_t &= a_0 + a_1 x + a_2 x^2 + a_3 x y + a_4 y; \\ y'_t &= b_0 + b_1 y + b_2 y^2 + b_3 x y + b_4 x, \\ a_j, b_j &\in \mathbb{R}. \end{aligned}$$

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Coupled Riccati equations (2)

$$\begin{aligned} x'_t &= a_0 + a_1 x + a_2 x^2 + a_3 x y + a_4 y; \\ y'_t &= b_0 + b_1 y + b_2 y^2 + b_3 x y + b_4 x, \\ a_j, b_j &\in \mathbb{R}. \end{aligned}$$

- Standard diffusive coupling terms a_4y, b_4x ;
- Multiplicative coupling terms *a*₃*xy*, *b*₃*xy*;
- Conditions for existence of soliton solutions ?
- Construction of soliton solutions ?

Inverse balancing technique

- The inverse idea of common ansatz methods;
- Insert known solution into DE;
- If DE depends linearly on equation parameters, solve for them;
- Application of technique is **not used to solve the DE**, but to obtain **necessary existence conditions** for the solitary solutions.

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PDEs with polynomial nonlinearity

Class of PDEs

$$\frac{\partial^m u}{\partial t^m} + A_{m-1,0} \frac{\partial^{m-1} u}{\partial t^{m-1}} + A_{0,m-1} \frac{\partial^{m-1} u}{\partial z^{m-1}} + \dots + A_{10} \frac{\partial u}{\partial t} + A_{01} \frac{\partial u}{\partial z} = a_n u^n + \dots + a_0.$$

When do the considered PDEs have solitary solutions (of any order l)?

$$u(t - \alpha z) = \frac{\sum_{k=0}^{l} \alpha_k \exp\left(\eta k \left(t - \alpha z\right)\right)}{\prod_{k=1}^{l} \left(\exp\left(\eta \left(t - \alpha z\right)\right) - t_k\right)}.$$

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Necessary existence conditions (1)

Condition #1: derivative and nonlinear term balance

n = m + 1.

Condition #2: equation and solution order balance

$$\frac{(m+1)l}{2} \le l+m+1, \qquad l,m \in \mathbb{N}.$$

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Necessary existence conditions (2)

Table of necessary existence conditions of solitary solutions to the considered PDEs. \exists denotes existence with all parameter values, \exists^* denotes existence with additional constraints on parameters, $\not\exists$ denotes the nonexistence of solitary solutions.

l (n,m)	(2,1)	(3, 2)	(4, 3)	(5, 4)	(6, 5)	(7, 6)	(8,7)
1	Ξ	∃*	∃*]*	∃*]*	∃*
2	A]*		*	3*	*	3*
3	A]*	*	*	3*	A	A
4	A	A	A	A	A	A	A
5	A	A	A	A	A	A	A

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Generalized differential operator

The notation
$$\mathbf{D}_{lpha} \coloneqq rac{\partial}{\partial lpha}$$
 will be used.

Generalized differential operator (GDO)

$$\mathbf{D}_{csu} \coloneqq R\left(c, s, u\right) \mathbf{D}_{c} + P\left(c, s, u\right) \mathbf{D}_{s} + Q\left(c, s, u\right) \mathbf{D}_{u},$$

where R, P, Q are analytic.

Properties of GDO

•
$$\mathbf{D}_{csu} \left(f_1 + f_2 \right) = \mathbf{D}_{csu} f_1 + \mathbf{D}_{csu} f_2;$$

•
$$\mathbf{D}_{csu}(f_1f_2) = (\mathbf{D}_{csu}f_1)f_2 + f_1\mathbf{D}_{csu}f_2.$$

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Multiplicative operator

Suppose a GDO D_{csu} is given. The multiplicative operator reads:

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$$\mathbf{G} := \sum_{j=0}^{+\infty} \frac{(t-c)^j}{j!} \mathbf{D}_{csu}^j.$$

Main property

$$\mathbf{G}f\left(c,s,u\right)=f\left(\mathbf{G}c,\mathbf{G}s,\mathbf{G}u\right).$$

Each ODE or ODE system has unique generalized differential and multiplicative operators.

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First order ODE system

$$\begin{split} & x'_t = P\left(t, x, y\right); \quad x = x\left(t; c, s, u\right), \quad x\left(c; c, s, u\right) = s; \\ & y'_t = Q\left(t, x, y\right); \quad y = y\left(t; c, s, u\right), \quad y\left(c; c, s, u\right) = u. \end{split}$$

System operators

$$\mathbf{D}_{csu} = \mathbf{D}_{c} + P\left(c, s, u\right) \mathbf{D}_{s} + Q\left(c, s, u\right) \mathbf{D}_{u}; \quad \mathbf{G} = \sum_{j=0}^{+\infty} \frac{(t-c)^{j}}{j!} \mathbf{D}_{csu}^{j}.$$

General solution

$$x(t) = \sum_{j=0}^{+\infty} \frac{(t-c)^j}{j!} \left(\mathbf{D}_{csu}^j s \right) = \mathbf{G}s, \quad y(t) = \sum_{j=0}^{+\infty} \frac{(t-c)^j}{j!} \left(\mathbf{D}_{csu}^j u \right) = \mathbf{G}u.$$

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Image of solution

Suppose an ODE is given:

$$x_{t}^{\prime}=P\left(x\right),\quad x=x\left(t;c,s\right);\quad x\left(c;c,s\right)=s.$$

Variable substitution

$$\widehat{t} := \exp(\eta t), \quad \widehat{c} := \exp(\eta c); \ \eta \in \mathbb{R} \setminus \{0\}$$

Image of solution

$$x = x(t) = x\left(\frac{1}{\eta}\ln\hat{t}\right) =: \hat{x}\left(\hat{t}\right) = \hat{x};$$
$$x'_t = \eta\hat{t}\hat{x}'_t.$$

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Image of ODE

Image of ODE

$$\begin{split} & \eta \widehat{t} \widehat{x}'_{\widehat{t}} = P\left(\widehat{x}\right); \\ & \widehat{x} = \widehat{x}\left(\widehat{t}; \widehat{c}, s\right), \quad \widehat{x}\left(\widehat{c}; \widehat{c}, s\right) = s. \end{split}$$

Operators of transformed ODE

$$\widehat{\mathbf{D}}_{\widehat{c}s} := \mathbf{D}_{\widehat{c}} + \frac{1}{\eta \widehat{c}} P(s) \mathbf{D}_{s};$$
$$\widehat{\mathbf{G}} := \sum_{j=0}^{+\infty} \frac{\left(\widehat{t} - \widehat{c}\right)^{j}}{j!} \widehat{\mathbf{D}}_{\widehat{c}s}^{j}.$$

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Linear recurring sequences

Let $p_j := \mathbf{D}_{csu}^j s$. Solutions can be written in the closed form if the sequence $(p_j; j \in \mathbb{Z}_0)$ (or a sequence constructed from p_j in a known way) is linearly recurring.

$$d_k := \det \left(\begin{bmatrix} p_0 & p_1 & \dots & p_{k-1} \\ p_1 & p_2 & \dots & p_k \\ \vdots & \vdots & \ddots & \vdots \\ p_{k-1} & p_k & \dots & p_{2k-2} \end{bmatrix} \right)$$

Linear recurring sequence

 $(p_j; j \in \mathbb{Z}_0)$ is an *m*-th order linear recurring sequence (LRS), if

$$d_m \neq 0, \quad d_{m+l} = 0; \quad l = 1, 2, \dots$$

It is denoted as

order
$$(p_j; j \in \mathbb{Z}_0) = m.$$

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Canonical expression of LRS

The characteristic equation reads:

$$\begin{vmatrix} p_0 & p_1 & \dots & p_{m-1} & p_m \\ p_1 & p_2 & \dots & p_m & p_{m+1} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ p_{m-1} & p_m & \dots & p_{2m-2} & p_{2m-1} \\ 1 & \rho & \dots & \rho^{m-1} & \rho^m \end{vmatrix} = 0.$$

Characteristic roots: $\rho_1, \ldots, \rho_m; \quad \rho_k \neq \rho_j, \, k \neq j.$

Canonical expression

$$p_j = \sum_{k=1}^m \lambda_k \rho_k^j, \quad j = 0, 1, \dots.$$

Coefficients $\lambda_k, k = 1, \dots, m$ are determined from a system of linear equations.

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Exponent sum solution

Suppose that

order
$$\left(\mathbf{D}_{csu}^{j}s; j \in \mathbb{Z}_{0}\right) = m;$$

 $\mathbf{D}_{csu}^{j}s = \sum_{k=1}^{m} \lambda_{k}\rho_{k}^{j}.$

Form of solution

$$x(t;c,s,u) = \mathbf{G}s = \sum_{j=0}^{+\infty} \frac{(t-c)^j}{j!} \sum_{k=1}^m \lambda_k \rho_k^j$$
$$= \sum_{k=1}^m \lambda_k \exp\left(\rho_k \left(t-c\right)\right).$$

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Solitary (soliton) solution (1)

Let

order
$$\left(\mathbf{D}_{csu}^{j}s; j \in \mathbb{Z}_{0}\right) = +\infty$$
, but order $\left(\frac{1}{j!}\widehat{\mathbf{D}}_{\widehat{c}su}^{j}s; j \in \mathbb{Z}_{0}\right) = m$.
Then $\frac{1}{j!}\widehat{\mathbf{D}}_{\widehat{c}su}^{j}s = \sum_{k=1}^{m} \lambda_{k}\rho_{k}^{j}$.

Theorem

The soliton solution exists, if the following relations hold true:

$$\widehat{\mathbf{D}}_{\widehat{c}su}\rho_k = \rho_k^2, \quad \widehat{\mathbf{D}}_{\widehat{c}su}\lambda_k = \lambda_k\rho_k; \quad k = 1, \dots, m.$$

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Solitary (soliton) solution (2)

Form of solution

$$\widehat{x}\left(\widehat{t};\widehat{c},s,u\right) = \widehat{\mathbf{G}}s = \sum_{j=0}^{+\infty} \left(\widehat{t}-\widehat{c}\right)^j \sum_{k=1}^m \lambda_k \rho_k^j = \sum_{k=1}^m \frac{\lambda_k}{1-\rho_k\left(\widehat{t}-\widehat{c}\right)};$$
$$x\left(t;c,s,u\right) = \frac{\sum_{k=0}^m \alpha_k \exp\left(\eta k(t-c)\right)}{\prod_{k=1}^m \left(\exp\left(\eta(t-c)\right)-t_k\right)}.$$

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General case

Suppose that
$$f(\xi) = \sum_{j=0}^{+\infty} \frac{q_j}{j!} \xi^j; \quad q_0 = 1, \, q_j = \prod_{k=0}^{j-1} (a+bk) \neq 0, j = 1, 2, \dots$$

Let order
$$\left(\frac{1}{q_j}\widehat{\mathbf{D}}_{\widehat{c}su}^j s; j \in \mathbb{Z}_0\right) = m, \quad \frac{1}{q_j}\widehat{\mathbf{D}}_{\widehat{c}su}^j s = \sum_{k=1}\lambda_k \rho_k^j.$$

Form of solution

$$\widehat{y}\left(\widehat{t};\widehat{c},s,u\right) = \widehat{\mathbf{G}}s = \sum_{j=0}^{+\infty} \frac{q_j}{j!} \left(\widehat{t}-\widehat{c}\right)^j \sum_{k=1}^m \lambda_k \rho_k^j = \sum_{k=1}^m \lambda_k f\left(\rho_k\left(\widehat{t}-\widehat{c}\right)\right);$$

$$x(t;c,s,u) = \sum_{k=1}^{m} \lambda_k f\left(\alpha_k - \beta_k \exp\left(\eta(t-c)\right)\right).$$

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Hepatitis C model (Coupled Riccati equations)

$$\begin{aligned} x'_t &= a_0 + a_1 x + a_2 x^2 + a_3 x y + a_4 y; \quad x(c) = s; \\ y'_t &= b_0 + b_1 y + b_2 y^2 + b_3 x y + b_4 x; \quad y(c) = u, \\ a_j, b_j &\in \mathbb{R}. \end{aligned}$$

- Kink solutions (order = 1);
- Bright/dark and singular solutions (order = 2);

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Bright/dark solutions

Analytical expression

$$x(t) = \sigma \frac{\left(\exp\left(\eta(t-c)\right) - x_1\right) \left(\exp\left(\eta(t-c)\right) - x_2\right)}{\left(\exp\left(\eta(t-c)\right) - t_1\right) \left(\exp\left(\eta(t-c)\right) - t_2\right)};$$

$$y(t) = \gamma \frac{\left(\exp\left(\eta(t-c)\right) - y_1\right) \left(\exp\left(\eta(t-c)\right) - y_2\right)}{\left(\exp\left(\eta(t-c)\right) - t_1\right) \left(\exp\left(\eta(t-c)\right) - t_2\right)}.$$
(1)

- σ, γ, η are constants;
- $x_1, x_2, y_1, y_2, t_1, t_2$ depend on initial conditions s, u;
- Solutions which hold for all initial conditions are constructed.

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Existence conditions

Bright/dark solitary solutions exist if:

$$a_3 = b_2; \quad a_2 = b_3,$$

and

$$9a_0a_1a_2 + 9b_0b_1b_2 - 18a_0a_2b_1 - 18b_0b_2a_1 + 3a_1b_1^2 + 3b_1a_1^2 - 2a_1^3 - 2b_1^3 - 9a_1a_4b_4 - 9b_1b_4a_4 + 27a_0b_2b_4 + 27b_0a_2a_4 = 0.$$

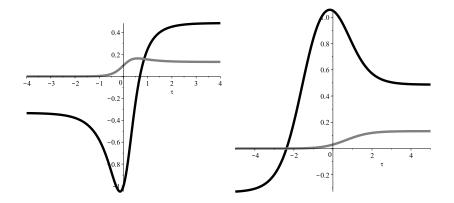
In the phase plane, bright/dark solution trajectories are conic sections:

$$Ax^2(t)+By^2(t)+Cx(t)y(t)+Ex(t)+Fy(t)=G;\quad A,B,C,E,F,G\in\mathbb{R}.$$

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Time evolution of solutions (1)



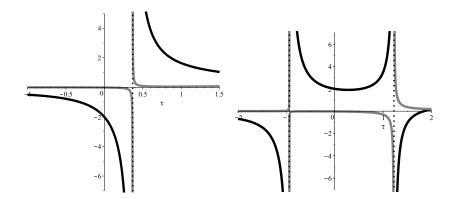
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Time evolution of solutions (2)

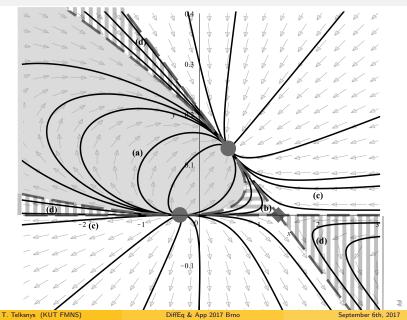


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Phase portrait



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PUBLICATIONS

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Publications I

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Publications

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Thank You For Your Attention

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